

## The Jost solutions to the Schrödinger equation with an additional complex potential

Khatira E. Abbasova, Agil Kh. Khanmamedov\*, Sevindj M. Bagirova

Received: 18.01.2022 / Revised: 27.04.2022 / Accepted: 09.05.2022

**Abstract.** We consider the differential equation  $-y'' + xe^{ix}y + q(x)y = k^2y$ . Using transformation operators, we obtain representations of solutions of this equation with conditions at infinity. Estimates for the kernels of the transformation operators are obtained.

**Keywords.** Schrödinger equation · non-self-adjoint differential operator · the space  $L_2(-\infty, +\infty)$  · transformation operator · the Jost solution.

**Mathematics Subject Classification (2010):** 34A55

### 1 Introduction and main results

In many aspects of the theory of inverse problems of spectral analysis, an important role is played by so-called transformation operators. The latter first appeared in the theory of generalized translation operators of J. Delsarte [1] and B.M. Levitan [5]. For arbitrary Sturm-Liouville equations, transformation operators were constructed by A.Ya. Povzner [9]. V.A. Marchenko [6] used transformation operators for studying inverse spectral problems and the asymptotic behavior of the spectral function of the singular Sturm-Liouville operator. It should be remarked that in the effective solution of various inverse problems of scattering theory, an important role is played by the transformation operators with a condition which were discovered by B.Ya. Levin [4]. Similar problems for the Schrödinger equation with unbounded potentials were considered in [3, 8, 10].

We consider the differential equation

$$-y'' + xe^{ix}y + q(x)y = k^2y, \quad (1.1)$$

---

\* Corresponding author

Kh.E. Abbasova,  
Azerbaijan State University of Economics (UNEC), AZ 1001, Baku, Azerbaijan  
E-mail: abbasova\_xatira@unec.edu.az

A.Kh. Khanmamedov  
Baku State University, AZ 1148, Baku, Azerbaijan  
Institute of Mathematics and Mechanics of NAS of Azerbaijan Baku, Azerbaijan  
Azerbaijan University, Baku, AZ-1073 Azerbaijan  
E-mail: agil.khanmamedov@yahoo.com

S.M. Bagirova  
Azerbaijan State Agricultural University, AZ 2000, Ganja, Azerbaijan  
E-mail: bagirovasevindj@rambler.ru

where  $q(x)$  is a continuously differentiable function with bounded support and is a complex parameter. If  $q(x) = 0$ , then from [2], the equation (1.1) has unique solution  $f_0(x, k)$ , which can be given as a series

$$f_0(x, k) = e^{ikx} + \sum_{n=1}^{\infty} \sum_{s=0}^n p_{ns}(k) x^s e^{i(n+k)x}. \quad (1.2)$$

Here  $p_{ns}(k)$  is a regular rational function with poles at the points  $k = -\frac{j}{2}$ ,  $j = 1, 2, \dots, n$  and multiplicities at most  $j + 1$ , while the series (1.2) admits term-by-term differentiation with respect to  $x$  any number of times for  $k \neq -\frac{n}{2}$ ,  $n = 1, 2, \dots$ . It was proved in [2] (see also [7]), for any  $k$  with  $\text{Im}k > 0$ , the function  $f_0(x, k)$  belongs to  $L_2(0, +\infty)$  and the function belongs to  $L_2(-\infty, 0)$ . Moreover, the functions  $f_0(x, k)$  and  $f_0(x, -k)$  form the fundamental system of solutions of equation (1.1) for  $k \neq 0$  when  $q(x) = 0$ .

This paper is devoted to the study of the solutions of (1.1) with asymptotic conditions

$$f_{\pm}(x, k) = f_0(x, \pm k) + o(1), \quad x \rightarrow \pm\infty.$$

We shall derive the integral representation, which is usually called the Jost translation representation between  $f_{\pm}(x, k)$  and  $f_0(x, \pm k)$ . The obtained results can be used to study the spectral properties of the non-self-adjoint differential operator  $L$ , generated by the differential expression  $l(y) = -y'' + xe^{ix}y + q(x)y$  in the space  $L_2(-\infty, +\infty)$ .

The main result of the present paper is as follows.

**Theorem 1.1** *For any  $k \neq -\frac{n}{2}$ ,  $n = 1, 2, \dots$  from the complex plane, equation (1.1) has solutions  $f_+(x, k)$  and  $f_-(x, k)$ , which can be represented in the form*

$$f_+(x, k) = f_0(x, k) + \int_x^{+\infty} K(x, t) f_0(t, k) dt \quad (1.3)$$

and

$$f_-(x, k) = f_0(x, -k) + \int_{-\infty}^x A(x, t) f_0(t, -k) dt. \quad (1.4)$$

Moreover,

$$K(x, x) = \frac{1}{2} \int_x^{+\infty} q(t) dt, \quad (1.5)$$

$$A(x, x) = \frac{1}{2} \int_{-\infty}^x q(t) dt. \quad (1.6)$$

## 2 Proof of the theorem

Without loss of generality, we consider the case " + " and assume that  $x \geq 0$ . We shall use the following notation

$$p(x) = xe^{ix}, \quad \sigma(x) = \frac{1}{2} \int_x^{+\infty} |q(t)| dt.$$

We first consider the following lemmas before turning to the proof of the theorem.

**Lemma 2.1** *If  $q(x)$  is a continuously differentiable function with bounded support, then the integral equation*

$$U(\xi_0, \eta_0) = \frac{1}{2} \int_{\xi_0}^{+\infty} q(\xi) d\xi + \int_0^{\eta_0} \int_{\xi_0}^{+\infty} [p(\xi - \eta) - p(\xi + \eta) + q(\xi - \eta)] U(\xi, \eta) d\xi d\eta \quad (2.1)$$

*has one and only one solution  $U(\xi_0, \eta_0)$ . Furthermore, if  $q(x) = 0$  when  $x > a$ , then*

$$U(\xi_0, \eta_0) = 0 \text{ when } \xi_0 \geq a. \quad (2.2)$$

**Proof.** Using the method of successive approximation, let

$$U_0(\xi_0, \eta_0) = \frac{1}{2} \int_{\xi_0}^{+\infty} q(\xi) d\xi, \quad (2.3)$$

$$U_n(\xi_0, \eta_0) = \int_0^{\eta_0} \int_{\xi_0}^{+\infty} [p(\xi - \eta) - p(\xi + \eta) + q(\xi - \eta)] U_{n-1}(\xi, \eta) d\xi d\eta. \quad (2.4)$$

Because the function  $q(x)$  with bounded support, there exists an  $a > 0$  such that  $q(x) = 0$  for  $x > a$ . By induction with respect to  $n$ , we have

$$U_n(\xi_0, \eta_0) \text{ for } \xi_0 > 2a, n = 0, 1, 2, \dots \quad (2.5)$$

For any  $R > 0$ , suppose that  $0 < \eta_0 < R$ ,  $0 < \xi_0 < +\infty$ . By (2.3), we have

$$|U(\xi_0, \eta_0)| \leq \sigma(\xi_0).$$

Taking the notation

$$M = \max_{\substack{0 \leq \xi \leq 2a \\ 0 \leq \eta \leq R}} |p(\xi - \eta) - p(\xi + \eta) + q(\xi - \eta)|$$

into account, we obtain

$$|U_1(\xi_0, \eta_0)| \leq \sigma(\xi_0) (M\eta_0).$$

Using induction, by (2.4) we next prove that

$$|U_n(\xi_0, \eta_0)| \leq \sigma(\xi_0) \frac{1}{n!} (M\eta_0)^n. \quad (2.6)$$

Hence the series

$$U(\xi_0, \eta_0) = \sum_{n=0}^{\infty} U_n(\xi_0, \eta_0) \quad (2.7)$$

is uniformly and absolutely convergent, so  $U(\xi_0, \eta_0)$  is the solution of the integral equation (2.1). From (2.6) and (2.7), it follows that

$$|U(\xi_0, \eta_0)| \leq \sigma(\xi_0) \exp(M\eta_0). \quad (2.8)$$

This implies obviously the uniqueness of the solution to the equation (2.1). The assertion (2.2) is justified by (2.5) and (2.7).

**Lemma 2.2** Suppose  $q(x)$  is a continuously differentiable function with bounded support. Then the solution  $U(\xi_0, \eta_0)$  of the integral equation (2.1) satisfies the following differential equation

$$\frac{\partial^2 U(\xi_0, \eta_0)}{\partial \xi_0 \partial \eta_0} + [p(\xi - \eta) - p(\xi + \eta) + q(\xi - \eta)] U(\xi_0, \eta_0) = 0 \quad (2.9)$$

and

$$U(\xi_0, 0) = \frac{1}{2} \int_{\xi_0}^{+\infty} q(\xi) d\xi. \quad (2.10)$$

**Proof.** From (2.1) the differentiability of  $U(\xi_0, \eta_0)$  is evident. Differentiating equation (2.1) directly, we get the equation (2.9). Putting  $\xi_0 = 0$  in (2.1), we get the result (2.10). We now let  $\xi_0 = \frac{t+x}{2}$ ,  $\eta_0 = \frac{t-x}{2}$  and express the function  $K(x, t) = U(\xi_0, \eta_0)$  as a function of  $x, t$ . Then the function  $K(x, t)$  is twice continuously differentiable. Moreover, from the two preceding lemmas we get the following lemma.

**Lemma 2.3** Suppose  $q(x)$  is a continuously differentiable function with bounded support. Then the function  $K(x, t) = U(\frac{t+x}{2}, \frac{t-x}{2})$  satisfies both the differential equation

$$\frac{\partial^2 K(x, t)}{\partial x^2} - [p(x) + q(x)] K(x, t) = \frac{\partial^2 K(x, t)}{\partial t^2} - p(t) K(x, t) \quad (2.11)$$

and the condition

$$K(x, x) = \frac{1}{2} \int_x^{+\infty} q(t) dt.$$

Furthermore, if  $q(x) = 0$  when  $x > a$ , then  $K(x, t) = 0$  when  $x + t > 2a$ .

Now the theorem can be proved. By differentiation from (1.3), we have

$$f'_+(x, k) = f'_0(x, k) - K'(x, x) f_0(x, k) + \int_x^{+\infty} K'_x(x, t) f_0(t, k) dt \quad (2.12)$$

$$\begin{aligned} f''_+(x, k) &= f''_0(x, k) - \frac{dK(x, x)}{dx} f_0(x, k) - K(x, x) f'_0(x, k) \\ &\quad - K'_x(x, t) f_0(x, k) + \int_x^{+\infty} K''_{xx}(x, t) f_0(t, k) dt. \end{aligned} \quad (2.13)$$

From Lemma 2.3, it is easily seen that when  $t$  sufficiently large,  $K(x, t) = 0$ , so the last terms of (1.3), (2.12), (2.13) are integrable. From

$$-f''_0(x, k) + p(x) f_0(x, k) = k^2 f_0(x, k) \quad (2.14)$$

and (1.3), we have

$$\begin{aligned} k^2 f_+(x, k) &= k^2 f_0(x, k) \\ &\quad + \int_x^{+\infty} K(x, t) p(t) f_0(t, k) dt - \int_x^{+\infty} K(x, t) f''_0(t, k) dt. \end{aligned} \quad (2.15)$$

Hence, integrating by parts, we obtain

$$\int_x^{+\infty} K(x, t) f''_0(t, k) dt = -K(x, x) f'_0(x, k) - \int_x^{+\infty} K'_t(x, t) f'_0(t, k) dt$$

$$= -K(x, x) f_0'(x, k) + K_t'(x, x) f_0(x, k) + \int_x^{+\infty} K_{tt}''(x, t) f_0(t, k) dt. \quad (2.16)$$

By virtue of (1.3) and (2.13)-(2.16), we have

$$\begin{aligned} & -f_+''(x, k) + p(x) f_+(x, k) - k^2 f_+(x, k) \\ &= \int_x^{+\infty} [K_{tt}''(x, t) - K_{xx}''(x, t) + K(x, t)(p(x) + q(x) - p(t))] f_0(t, k) dt \\ & \quad + \left[ 2 \frac{dK(x, x)}{dx} + q(x) \right] f_0(x, k). \end{aligned}$$

From the lemma 2.3 and the last relation,  $f_+(x, k)$  satisfies equation (1.1). Furthermore, by virtue of (2.8)-(2.14), it follows that  $f_+(x, k) = f_0(x, k)$  when  $x$  sufficiently large. Hence, the  $f_+(x, k)$  is a Jost solution. Thus, the proof of the theorem is complete.

## References

1. Delsarte J.: Sur une extension de la formule de Taylor, *J. Math. Pures Appl.* **17**, 213-231 (1938).
2. Gasymov M.G.: On the spectrum of a non-self-adjoint operator, *Uspekhi Mat. Nauk*, **36**(6), 209-210 (1981).
3. Gasymov M.G., Mustafayev B.A.: Inverse scattering problem for an anharmonic equation on the semiaxis, *Dokl. Akad. Nauk SSSR*, **228**(1), 11-14 (1976).
4. Levin B.Ya.: Transformation of Fourier and Laplace type by means of solutions of a differential equation of second order, *Dokl. Akad. Nauk SSSR*, **106**(2), 187-190 (1956).
5. Levitan B.M.: *Inverse SturmLiouville Problems* (Nauka, Moscow, 1984) [in Russian].
6. Marchenko V.A.: Some questions of the theory of differential operators of second order, *Dokl. Akad. Nauk SSSR*, **72**(3), 457-460 (1950).
7. Manafov M.Dzh. : Spectrum and Spectral Decomposition of a Non-Self-Adjoint Differential Operator, *Math. Notes*, **82**(1), 52-56 (2007.)
8. Masmaliev G.M., Khanmamedov A.Kh.: Transformation Operators for Perturbed Harmonic Oscillators, *Math. Notes*, **105**(5), 728-733(2019).
9. Povzner A.Ya.: On differential equations of Sturm-Liouville type on a half-axis, *Mat. Sb.* **23**(65) (1), 3-52 (1948).
10. Yishen Li.: One special inverse problem of the second-order differential equation for the whole real axis, *Chinese Ann. Math.*, **2** (2), 147-155 (1981).