

The mixed problem for nonlinear Timoshenko systems

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Abstract. *In this paper, a mixed problem is considered for a system of fourth order multidimensional semilinear hyperbolic equations with constant coefficients. Under suitable hypotheses we prove existence of a global solution.*

Keywords. system of fourth order wave equation · global solvability · Timoshenko system · mixed problem.

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1 Introduction

The linear theory of the vibration of the bar is described by a system of differential equations

$$\left. \begin{aligned} E I y_{xxxx} + \rho A y_{tt} - \rho A e \theta_{tt} &= g_1(t, x) \\ E C_w \theta_{xxxx} - G C \theta_{xx} - \rho A e y_{tt} + \rho (I + A e^2) \theta_{tt} &= g_2(t, x) \end{aligned} \right\}, \quad (1.1)$$

where $0 < x < l$, $0 < t < T$, $l > 0$, $T > 0$ are given numbers, $y(x, t)$ is a transverse displacement, $\theta(x, t)$ is an angle of cross-section of a bar, E is the Young modulus, I is a polar moment of inertia of the cross section with respect to its center of gravity, ρ is a density of the material of the bar, A is a cross-sectional area, e is a distance from center of gravity to center of torsion, C_w is a sectorial moment of inertia of the cross section, G is a shear modulus, C is a geometric rigidity of free torsion, EC_w is a stiffness of bending torsion, GC is a stiffness of free torsion [1,2].

2 Formulation of the problem and main results

In this work we study a mixed problem for a multidimensional semilinear analogue of the system (1.1). Let $\Omega \subset R^n$ be a bounded domain with smooth boundary Γ . In the domain

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$Q_T = [0, T] \times \Omega$ we consider the mixed problem:

$$\left. \begin{aligned} EI\Delta^2 y + \rho A y_{tt} - \rho A e \theta_{tt} &= f_1(t, x, y, \theta) \\ EC_w \Delta^2 \theta - GC \Delta \theta - \rho A e y_{tt} + \rho (I + Ae^2) \theta_{tt} &= f_2(t, x, y, \theta) \end{aligned} \right\} \quad (2.1)$$

with boundary conditions

$$\left. \begin{aligned} y(t, x) = \Delta y(t, x) &= 0, \quad t \in [0, T], \quad x \in \Gamma \\ \theta(t, x) = \Delta \theta(t, x) &= 0, \quad t \in [0, T], \quad x \in \Gamma \end{aligned} \right\} \quad (2.2)$$

and initial conditions

$$\left. \begin{aligned} y(0, x) = y_0(x), \quad y_t(0, x) = y_1(x) \\ \theta(0, x) = \theta_0(x), \quad \theta_t(0, x) = \theta_1(x) \end{aligned} \right\}, \quad x \in \Omega, \quad (2.3)$$

where $t \in [0, T]$, $x \in \Omega$, Δ is a Laplace operator. We assume that all the coefficients A, C, E, G, I, C_w, e and ρ are positive constants. We denote by $\langle w^1, w^2 \rangle$ the scalar product in $L_2(0, 1)$. We will investigate the problem (2.1) - (2.3) in the Hilbert space $\mathcal{H} = L_2(0, 1) \times L_2(0, 1)$ with the scalar product of \mathcal{H}

$$\langle w^1, w^2 \rangle = \frac{I}{C_w} \langle y^1, y^2 \rangle_{L_2(0,1)} + \langle \theta^1, \theta^2 \rangle_{L_2(0,1)},$$

where $w^k = \begin{pmatrix} y_k \\ \theta_k \end{pmatrix}$, $k = 1, 2$.

Let us define the functional spaces

$$\widehat{H}_0^{2s} = \left\{ u : u \in H^{2s} = W_2^{2s}(\Omega), \Delta^l u(x) = 0, x \in \Gamma, l = 0, 1, \dots, s-1 \right\}.$$

We introduce the following notation: $\mathcal{H}_1 = \widehat{H}_0^2 \times \widehat{H}_0^2$, $\mathcal{H}_2 = \widehat{H}_0^4 \times \widehat{H}_0^4$.

Let L be a linear operator defined in the space, where

$$D(L) = \mathcal{H}_2,$$

$$Lw = \begin{bmatrix} \frac{E(I+Ae^2)}{\rho A} \Delta^2 & \frac{eEC_w}{\rho I} \Delta^2 \\ \frac{eE}{\rho} \Delta^2 & \frac{EC_w}{\rho I} \Delta^2 \end{bmatrix} w, \quad w = \begin{pmatrix} y \\ \theta \end{pmatrix} \in D(L).$$

We also define the linear operator L_1 as follows

$$D(L_1) = \mathcal{H}_1,$$

$$L_1 w = \begin{bmatrix} 0 - \frac{eGC}{\rho I} \Delta \\ 0 - \frac{GC}{\rho I} \Delta \end{bmatrix} w, \quad w = \begin{pmatrix} y \\ \theta \end{pmatrix} \in D(L_1).$$

Then problem (2.1), (2.3) can be written in the form

$$\left. \begin{aligned} w_{tt} + Lw + L_1 w &= G(t, w) \\ w(0) = w_0, \quad w_t(0) &= w_1 \end{aligned} \right\}, \quad (2.4)$$

where $w(t) = \begin{pmatrix} y(t, x) \\ \theta(t, x) \end{pmatrix}$, $w_0 = \begin{pmatrix} y_0(x) \\ \theta_0(x) \end{pmatrix}$, $w_1 = \begin{pmatrix} y_1(x) \\ \theta_1(x) \end{pmatrix}$. Here $G(t, x)$ is a nonlinear operator defined as follows

$$G(t, w) = \begin{pmatrix} h_1(t, x, w) \\ h_2(t, x, w) \end{pmatrix},$$

where

$$h_1(t, x, w) = \frac{I + Ae^2}{\rho AI} f_1(t, x, y, \theta) + \frac{e}{\rho I} f_2(t, x, y, \theta)$$

$$h_2(t, x, w) = \frac{e}{\rho I} f_1(t, x, y, \theta) + \frac{1}{\rho I} f_2(t, x, y, \theta).$$

Suppose that the following conditions are satisfied:

- i) $f_k(t, x, y, \theta) \in C^1([0, T] \times \Omega \times R^2)$, $k = 1, 2$;
- ii) if $n \geq 4$, then

$$|f_k(t, x, y, \theta)| \leq c(1 + |y|^p + |\theta|^q),$$

$$|f'_{k_t}(t, x, y, \theta)| \leq c(1 + |y|^p + |\theta|^q),$$

$$|f'_{k_y}(t, x, y, \theta)| \leq c(1 + |y|^{p-1} + |\theta|^{q-1}),$$

$$|f'_{k_\theta}(t, x, y, \theta)| \leq c(1 + |y|^{p-1} + |\theta|^{q-1}),$$

where $p, q \geq 1$ and $p, q \leq \frac{n}{n-4}$ if $n > 4$.

Using the definitions of L, L_1 and $G(t, w)$ the following lemmas are proved.

Lemma 2.1 L is a positive self-adjoint operator in \mathcal{H} .

Lemma 2.2 The linear operator L_1 is subjected to the operator $L^{\frac{1}{2}}$, i.e. $D(L^{\frac{1}{2}}) \subset D(L)$ and $\|L_1 w\|_{\mathcal{H}}^2 \leq c \langle Lw, w \rangle$, $c > 0$, $w \in D(L^{\frac{1}{2}})$,

Lemma 2.3 Let the conditions i) and ii) be satisfied. Then $G(t, w)$ acts from \mathcal{H}_1 to \mathcal{H} and satisfies the local Lipschits condition, i.e. for any $t_1, t_2 \in [0, T]$ and $w^1, w^2 \in \mathcal{H}_1$ the following inequalities hold

$$\|G(t_1, w^1) - G(t_2, w^2)\|_{\mathcal{H}} \leq c \left(\|w^1\|_{\mathcal{H}_1}, \|w^2\|_{\mathcal{H}_1} \right) \times \left[|t_1 - t_2| + \|w^1 - w^2\|_{\mathcal{H}_1} \right],$$

where $c(\cdot) \in C(R_+^2, R_+)$.

From Lemmas 2.1-2.3 and [3] theorem on local solvability follows:

Theorem 2.1 Let the conditions i) and ii) be satisfied. Then for any $(y_0, \theta_0) \in \mathcal{H}_1, (y_1, \theta_1) \in \mathcal{H}$ there is a $T' > 0$, such that the problem (2.1)-(2.3) has a unique solution (y, θ) , where

$$(y, \theta) \in C([0, T'], \mathcal{H}_1 \times \mathcal{H}_1) \cap C([0, T'], \mathcal{H} \times \mathcal{H}),$$

Moreover, if T_{max} is the length of the maximum interval of the existence of solutions, then one of the following alternatives is fulfilled:

$$i) \lim_{t \rightarrow T_{max} - 0} \left[\|y_t(t, \cdot)\|_{L_2(\Omega)}^2 + \|\theta_t(t, \cdot)\|_{L_2(\Omega)}^2 + \|y(t, \cdot)\|_{H^1}^2 + \|\theta(t, \cdot)\|_{H^1}^2 \right] = +\infty$$

or

ii) $T_{max} = T$.

Now consider the mixed problem

$$\left. \begin{aligned} \rho A y_{tt} - \rho A e \theta_{tt} + EI \Delta^2 y + |y|^p |\theta|^{p+2} y &= g_1(t, x) \\ -\rho A e y_{tt} + \rho (I + Ae^2) \theta_{tt} + EC_w \Delta^2 \theta - GC \Delta \theta + |y|^{p+2} |\theta|^p \theta &= g_2(t, x) \end{aligned} \right\} \quad (2.5)$$

with boundary and initial conditions (2.2),(2.3).

We will suppose that, the following conditions are satisfied

- 1) $g_k(t, x) \in L_2((0, T); L_2(\Omega))$, $k = 1, 2$,
- 2) $p > 0$ and $0 < p < \frac{8}{n-4}$ if $n > 4$.

Theorem 2.2 *Let the conditions 1) and 2) be satisfied. Then for any*

$$y_0, \theta_0 \in \widehat{H}_0^2, \quad y_1, \theta_1 \in L_2(\Omega).$$

There exist functions $y, \theta : (0, T) \rightarrow L_2(\Omega)$ such that

$$\begin{aligned} y, \theta &\in L^\infty([0, T], \widehat{H}_0^2), \\ y', \theta' &\in L^\infty([0, T], L_2(\Omega)), \\ y\theta &\in L^\infty([0, T], L^{p+2}(\Omega)), \end{aligned}$$

satisfying the nonlinear systems

$$\rho A y_{tt} - \rho A e \theta_{tt} + EI \Delta^2 y + |y|^p |\theta|^{p+2} y = g_1(t, x)$$

in $L_2(0, T; H^{-2}(\Omega)) \cap L^q(Q)$,

$$-\rho A e y_{tt} + \rho (I + A e^2) \theta_{tt} + EC_w \Delta^2 \theta - GC \Delta \theta + |y|^{p+2} |\theta|^p \theta = g_2(t, x)$$

in $L_2(0, T; H^{-2}(\Omega)) \cap L^q(Q)$.

3 Proof of Theorem 2.2.

We give the proof of Theorem 2.2 for $n > 4$. In the case $n \leq 4$, Theorem 2.2 follows from Theorem 2.1 and the corresponding a priori estimate.

We define the numbers $q > 1$ and $q' > 1$ in the following way:

$$q = \frac{2n}{(n-4)(p+2) + 2n(p+1)}, \quad q' = \frac{2n(p+2)}{(n+4)(p+2) - 2n(p+1)}$$

It is obvious that $\frac{1}{q} + \frac{1}{q'} = 1$. We choose $\rho > 1$ and $\rho' > 1$ in the following way

$$\rho = \frac{p+2}{p+1} \cdot \frac{1}{q}, \quad \rho' = \frac{\rho}{\rho-1}.$$

Using the Holder inequality, we have

$$\int_{\Omega} (|y|^{p+2} |\theta|^{p+1})^q dx \leq \left(\int_{\Omega} |y\theta|^{(p+1)qp} dx \right)^{1/p} \cdot \left(\int_{\Omega} |y|^{q\rho'} dx \right)^{1/\rho'}$$

Taking into account (2.5) and embedding theorems, we obtain that

$$\begin{aligned} \int_{\Omega} (|y|^{p+2} |\theta|^{p+1})^q dx &\leq \|y\theta\|_{L_{p+2}(\Omega)}^{(p+1)q} \cdot \|y\|_{L_{\frac{2n}{n-4}}}^q \leq \\ &\leq c \|y\theta\|_{L_{p+2}(\Omega)}^{(p+1)q} \cdot \|y\|_{L_{\widehat{H}_0^2}}^q. \end{aligned} \quad (3.1)$$

We will solve the problem (2.5), (3.1) by the Galerkin method :

Let $\{w_k\}$ be a basis in $\widehat{H}_0^2 \cap L^{q_1}(\Omega)$, where $q_1 = \max\{p+2, q'\}$. We denote by V_m the subspace spanned by the first m vectors $\{w_1, w_2, \dots, w_m\}$

Let

$$y_m = \sum_{j=1}^m u_{j_m}(t) w_j,$$

$$\theta_m = \sum_{j=1}^m v_{j_m}(t) w_j,$$

where $(u_{j_m}(t), v_{j_m}(t))$ is defined from the following system of equations:

$$\begin{aligned} \rho A (y_m'', w_j) - \rho A e (\theta_m'', w_j) + EI (\Delta y, \Delta w_j) + (|y_m|^p |\theta_m|^{p+2} y_m, w_j) \\ = (g_1(t, x), w_j) \end{aligned} \quad (3.2)$$

$$\begin{aligned} -\rho A e (y_m'', w_j) + \rho (I + A e^2) (\theta_m'', w_j) + EC_w (\Delta \theta, \Delta w_j) + GC (\nabla_x \theta, \nabla w_j) \\ + (|y_m|^p |\theta_m|^{p+2} \theta_m, w_j) = (g_2(t, x), w_j). \end{aligned} \quad (3.3)$$

The system (3.2), (3.3) is considered with the initial conditions

$$\left. \begin{aligned} y_m(0) &= y_{0m}, & y_m'(0) &= y_{1m} \\ \theta_m(0) &= \theta_{0m}, & \theta_m'(0) &= \theta_{1m} \end{aligned} \right\}. \quad (3.4)$$

Here $y_{0m}, y_{1m}, \theta_{0m}, \theta_{1m} \in V_m$ and

$$\begin{aligned} y_{0m} &\rightarrow y_0 & \text{in } \widehat{H^2} \cap L^p(\Omega), \\ y_{1m} &\rightarrow y_1 & \text{in } L_2(\Omega), \\ \theta_{0m} &\rightarrow \theta_0 & \text{in } \widehat{H^2} \cap L^p(\Omega), \\ \theta_{1m} &\rightarrow \theta_1 & \text{in } L_2(\Omega). \end{aligned}$$

By the theorem on the existence of the solutions of the Cauchy problem for the system of differential equations, there exist numbers $T_m > 0$, such that the problem (3.2) - (3.4) has a solution

$$(u_{j_m}(t), v_{j_m}(t)) \in C^1[0, T_m], \quad j = 1, \dots, m, \quad m = 1, 2, \dots.$$

A priori estimates

We multiply the equation (3.2) by $y'_{j_m}(t)$ and sum from $j = 0$ to $j = m$. Then we have

$$\begin{aligned} \frac{\rho A}{2} \frac{d}{dt} \int_{\Omega} |y'_m(t, x)|^2 dx - \rho A e \int_{\Omega} \theta_m''(t, x) y'_m(t, x) dx + \frac{EI}{2} \frac{d}{dt} \int_{\Omega} |\Delta y_m(t, x)|^2 dx \\ + \int_{\Omega} |y_m(t, x)|^p |\theta_m(t, x)|^{p+2} y_m(t, x) y'_m(t, x) dx = \int_{\Omega} g_1(t, x) y'_m(t, x) dx. \end{aligned} \quad (3.5)$$

In the similar way, multiplying the equation (3.3) by $\theta'_{j_m}(t)$ and summing from $j = 0$ to $j = m$, we have :

$$\begin{aligned} -\rho A e \int_{\Omega} y_m''(t, x) \theta'_m(t, x) dx \\ + \frac{\rho (I + A e^2)}{2} \frac{d}{dt} \int_{\Omega} |\theta'_m(t, x)|^2 dx + \frac{EC_w}{2} \frac{d}{dt} \int_{\Omega} |\Delta \theta_m(t, x)|^2 dx \\ + \frac{GC}{2} \int_{\Omega} |\nabla \theta_m(t, x)|^2 dx + \int_{\Omega} |y_m(t, x)|^{p+2} |\theta_m(t, x)|^p \theta_m(t, x) \theta'_m(t, x) dx \end{aligned}$$

$$= \int_{\Omega} g_2(t, x) \theta'_m(t, x) dx. \quad (3.6)$$

Summing (3.5) and (3.6), we get

$$\begin{aligned} & \frac{d}{dt} \left[\frac{\rho A}{2} \int_{\Omega} |y'_m(t, x)|^2 dx \right. \\ & + \frac{EI}{2} \int_{\Omega} |\Delta y_m(t, x)|^2 dx + \frac{\rho(I + Ae^2)}{2} \int_{\Omega} |\theta'_m(t, x)|^2 dx \\ & \quad + \frac{EC_w}{2} \int_{\Omega} |\Delta \theta_m(t, x)|^2 dx \\ & \quad \left. - \frac{GC}{2} \int_{\Omega} |\nabla \theta_m(t, x)|^2 dx - \rho A \int_{\Omega} \theta'_m(t, x) y'_m(t, x) dx \right] \\ & \quad + \frac{1}{p+2} \frac{d}{dt} \int_{\Omega} |y_m(t, x)|^{p+2} |\theta_m(t, x)|^{p+2} dx \\ & = \int_{\Omega} f_1(t, x) y'_m(t, x) dx + \int_{\Omega} f_2(t, x) \theta'_m(t, x) dx. \end{aligned} \quad (3.7)$$

Using the Holder and Young inequality we have

$$\rho A e \int_{\Omega} \theta'_m(t, x) y'_m(t, x) dx \leq \frac{\rho A e}{2\varepsilon} \int_{\Omega} |\theta'_m(t, x)|^2 dx + \frac{\rho A}{2} \varepsilon \int_{\Omega} |y'_m(t, x)|^2 dx, \quad (3.8)$$

where $0 < \varepsilon < 1$ and

$$\frac{\rho A e}{2\varepsilon} < \frac{\rho(I + Ae^2)}{2}, \quad (3.9)$$

i.e.

$$0 < \varepsilon < \frac{Ae}{I + Ae^2}.$$

It is obvious that $0 < \frac{Ae}{I + Ae^2} < 1$.

Estimating the right-hand side from above, we obtain

$$\begin{aligned} & \int_{\Omega} f_1(t, x) y'_m(t, x) dx + \int_{\Omega} f_2(t, x) \theta'_m(t, x) dx \\ & \leq c + \int_{\Omega} |y'_m(t, x)|^2 dx + \int_{\Omega} |\theta'_m(t, x)|^2 dx. \end{aligned} \quad (3.10)$$

From (3.5)-(3.10), we obtain that

$$\begin{aligned} & \frac{\rho A}{2} (1 - \varepsilon) \int_{\Omega} |y'_m(t, x)|^2 dx + \left(\frac{\rho(I + Ae^2)}{2} - \frac{\rho A e}{2\varepsilon} \right) \int_{\Omega} |\theta'_m(t, x)|^2 dx \\ & + \frac{EI}{2} \int_{\Omega} |\Delta y_m(t, x)|^2 dx + \frac{1}{p+2} \int_{\Omega} |y_m(t, x) \theta_m(t, x)|^{p+2} dx \\ & \leq \frac{1}{2} \int_0^t \left[\int_{\Omega} |y'_m(s, x)|^2 dx + \int_{\Omega} |\theta'_m(s, x)|^2 dx \right] ds. \end{aligned}$$

Using Gronwall's Lemma, from here we get:

$$\int_{\Omega} |y'_m(t, x)|^2 dx \leq c, \quad (3.11)$$

$$\int_{\Omega} |\theta'_m(t, x)|^2 dx \leq c, \quad (3.12)$$

$$\int_{\Omega} |\Delta y_m(t, x)|^2 dx \leq c, \quad (3.13)$$

$$\int_{\Omega} |\Delta \theta_m(t, x)|^2 dx \leq c, \quad (3.14)$$

$$\int_{\Omega} |y_m(t, x) \theta_m(t, x)|^{p+2} dx \leq c. \quad (3.15)$$

On the other hand, using (3.11), (3.13) - (3.15), embedding theorem [4], we obtain that

$$\begin{aligned} & \int_{\Omega} \left(|y_m(t, x)|^{p+2} |\theta_m(t, x)|^{p+1} \right)^q dx \\ & \leq \|y_m(t, x) \theta_m(t, x)\|_{L_{p+2}(\Omega)}^{(p+1)q} \|y_m\|_{\widehat{H}_0^2}^q \leq c. \end{aligned} \quad (3.16)$$

Hence we have the following estimate

$$\int_{\Omega} \left(|y_m(t, x)|^{p+1} |\theta_m(t, x)|^{p+2} \right)^q dx \leq c. \quad (3.17)$$

In the similar way, we obtain

$$\int_{\Omega} \left(|y_m(t, x)|^{p+2} |\theta_m(t, x)|^{p+1} \right)^q dx \leq c. \quad (3.18)$$

By (3.11)-(3.15) from the sequence $\{(y_m, \theta_m)\}$ we can select a sequence $\{(y_\mu, \theta_\mu)\}$, such that

$$y_\mu \rightharpoonup y \text{ weakly star in } L_\infty(0, T; H_0^2), \quad (3.19)$$

$$\theta_\mu \rightharpoonup \theta \text{ weakly star in } L_\infty(0, T; H_0^2), \quad (3.20)$$

$$y'_\mu \rightharpoonup y' \text{ weakly star in } L_\infty(0, T; L_2(\Omega)), \quad (3.21)$$

$$\theta'_\mu \rightharpoonup \theta' \text{ weakly star in } L_\infty(0, T; L_2(\Omega)), \quad (3.22)$$

$$y_\mu \theta_\mu \rightharpoonup \chi \text{ weakly star in } L_\infty(0, T; L_{p+2}(\Omega)). \quad (3.23)$$

It follows from (3.19)-(3.22) that

$$y_\mu \longrightarrow y \text{ strongly in } L_2(Q_T), \quad (3.24)$$

$$\theta_\mu \longrightarrow \theta \text{ strongly in } L_2(Q_T), \quad (3.25)$$

$$y_\mu \longrightarrow y \text{ a.e. in } Q_T, \quad (3.26)$$

$$\theta_\mu \longrightarrow \theta \text{ a.e. in } Q_T. \quad (3.27)$$

It follows from (3.26), (3.27) that

$$|y_\mu|^p |\theta_\mu|^{p+2} y_\mu \longrightarrow |y|^p |\theta|^{p+2} y \text{ a.e. } Q_T \quad (3.28)$$

and

$$|y_\mu|^{p+2}|\theta_\mu|^p\theta_\mu \longrightarrow |y|^{p+2}|\theta|^p\theta \quad \text{a.e. } Q_T. \quad (3.29)$$

On the other hand, it follows from (3.17) that $\left\{ |y_\mu|^p|\theta_\mu|^{p+2}y_\mu \right\}$ is bounded in $L_\infty(0, T; L^q(\Omega))$. Then in view of (3.28)

$$|y_\mu|^{p+2}|\theta_\mu|^p y_\mu \rightharpoonup |y|^{p+2}|\theta|^p y \quad \text{weakly star in } L_\infty(0, T; L^q(\Omega)).$$

Similarly, it follows from (3.18) and (3.29) that

$$|y_\mu|^p|\theta_\mu|^{p+2}\theta_\mu \rightharpoonup |y|^p|\theta|^{p+2}\theta \quad \text{weakly star in } L_\infty(0, T; L^q(\Omega)).$$

It follows from (3.24), (3.25) that

$$y_\mu\theta_\mu \longrightarrow y\theta \quad \text{in } L^1(Q). \quad (3.30)$$

By (3.23), (3.30)

$$\chi = y\theta.$$

Passing to the limit is carried out as follows:

Let $\eta \in D(0, T)$. Multiplying both sides of (3.3) by η and integrating parts, we obtain that

$$\begin{aligned} & -\rho A \int_0^T \langle y'_m, \eta' w_j \rangle dt + \rho A e \int_0^T \langle \theta'_m, \eta' w_j \rangle dt + EI \int_0^T \langle \Delta y, \eta w_j \rangle dt + \\ & + \int_0^T \langle |y_m|^p |\theta_m|^{p+2} y_m, \eta w_j \rangle dt = \int_0^T \langle g_1, \eta w_j \rangle dt. \end{aligned}$$

Passing to the limit, hence we get :

$$\begin{aligned} & -\rho A \int_0^T \langle y', \eta' w_j \rangle dt + \rho A e \int_0^T \langle \theta', \eta' w_j \rangle dt + EI \int_0^T \langle \Delta y, \eta w_j \rangle dt \\ & + \int_0^T \langle |y|^p |\theta|^{p+2} y, \eta w_j \rangle dt = \int_0^T \langle g_1, \eta w_j \rangle dt. \end{aligned}$$

The set of functions $\{\eta w_1, \eta w_2, \dots, \eta w_n\}$ is dense in $H_0^1(0, T; H_0^1(\Omega)) \cap L^{p_1}(Q)$, therefore

$$\begin{aligned} & -\rho A \int_0^T \langle y', z' \rangle dt + \rho A e \int_0^T \langle \theta', \eta' w_j \rangle dt + EI \int_0^T \langle \Delta y_{xx}, \Delta z \rangle dt \\ & + \int_0^T \langle |y|^p |\theta|^{p+2} y, z \rangle dt = \int_0^T \langle g_1, z \rangle dt. \end{aligned}$$

It is true for any $z \in H_0^1(0, T; H_0^2(\Omega)) \cap L^{p_1}(Q)$.

Similarly, we get that, for any $z \in H_0^1(0, T; H_0^2(\Omega)) \cap L^{p_1}(Q)$

$$\begin{aligned} & -\rho A e \int_0^T \langle y', z' \rangle dt + \rho (I + A e^2) \int_0^T \langle \theta', z' \rangle dt + EC_w \int_0^T \langle \Delta \theta, \Delta z \rangle dt \\ & + GC \int_0^T \langle \Delta \theta, \nabla w_j \rangle dt + \int_0^T \langle |y|^{p+2} |\theta|^p \theta, z \rangle dt = \int_0^T \langle g_2, z \rangle dt. \end{aligned}$$

The uniqueness result via the standard energy method (see[5], pp.32-33).

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